

LPG Fiber Optic Sensor Applied to the Determination of the Flexural Elasticity Modulus of Woods

Luis Mosquera, L. Yana and Jesús Basurto

EasyChair preprints are intended for rapid dissemination of research results and are integrated with the rest of EasyChair.

February 10, 2020

# LPG FIBER OPTIC SENSOR APPLIED TO THE DETERMINATION OF THE FLEXURAL ELASTICITY MODULUS OF WOODS

#### L Mosquera<sup>1</sup>, L Yana<sup>1</sup>, J Basurto<sup>1</sup>

<sup>1</sup>Facultad de Ingeniería Civil, Universidad Nacional de Ingeniería, Apartado 1301, Lima Perú

# lmosquera@uni.edu.pe

## Abstract

In this work, we show the feasibility of using  $CO_2$  laser induced long-period gratings (LPGs) for determining the flexural elasticity modules of woods. To do this, we characterized the dynamic response of bars made of different woods put in oscillation. By recovering the bars' flexural oscillations from the LPG time response and by taking the fast Fourier transform (FFT) of it, the movement vibration frequencies could be obtained. Knowledge of these vibration frequencies for the oscillation lengths of different bars allowed us to calculate the materials' the flexural elasticity modules.

Keywords: optical fiber sensing; vibration sensor; MOE of woods

#### 1. The flexural elasticity modules determination

In the 'bar resonance' method, a bar of the material of interest is arranged in a cantilever configuration and put in oscillation. The resulting movement is studied in order to obtain the bar material's flexural elasticity modules value [1,3]. We use a curvature sensitive long-period fiber grating to register the displacement of the bar as a function of time.

If the bar in cantilever configuration is vertically deflected from its equilibrium state and then released, the resulting movement is oscillatory and its amplitude decays as a function of time. The formal treatment of this movement, via the Euler–Bernoulli equation, indicates that the general solution has a series form with infinite frequency components, which are given by equation (1), where  $\rho$  and Y are, respectively, the material density and flexural elasticity modules; A and L are, the cross-sectional area and the vibrating length of the bar; I is the second moment of the cross-section, which, for a rectangular bar of width b and depth h, can be calculated as b h<sup>3</sup>/12; kv is the viscous damping constant of the system. The value of the constant  $\lambda n$  is determined by the boundary conditions of the problem.

$$f_{n} = \frac{1}{2\pi} \sqrt{\frac{YI}{\rho A} \left(\frac{\lambda_{n}}{l}\right)^{4} - \left(\frac{k_{v}}{2\rho A}\right)^{2}}$$
(1)

As we focus on the determination of the Young's modulus, one can see, from equation (1), that the knowledge of only one frequency component is enough to attain our goal. In this investigation, we determined the first frequency component,  $f_1$ , and then found the flexural elasticity modules of the materials of interest.

## 2. Long-period gratings and their curvature sensitivity

Long-period fiber gratings (LPGs) consist of a longitudinal periodic perturbation of the refractive index of an optical fiber which is able to provide coupling between core and cladding modes at certain wavelengths. Equation (2) describes the wavelengths  $\lambda^{(m)}$  where the coupling between the referenced modes happen—nco is the effective refractive index of the core mode,  $n_{cl}^m$  is the effective refractive index of the mth order cladding mode and  $\Lambda$  is the period of the refractive index perturbation [4].

$$\lambda^{(m)} = \left( n_{co} - n_{cl}^{(m)} \right) \Lambda \tag{2}$$

The curvature causes stretching or compression of the fiber by changing the refractive indices  $n_{co}$ ,  $n_{cl}$  due to the elastooptic effect. In addition, the elastic deformation changes the period of the grating ( $\Lambda$ ). Changes in refractive indexes of the fiber and in the period of the LPG cause a shift in the wavelength of resonance  $\lambda$  as well as, changes in the depth of the transmittance dip (equation (3)).

$$T^{(m)} = 1 - \sin^2(\kappa^{(m)}L)$$
(3)

Where,  $T^{(m)}$  is the dip of the m-th mode transmittance and  $\kappa^{(m)}$  denotes coupling coefficient between the core mode and the m-th order cladding mode given as

$$\kappa^{(m)} = \frac{\pi \Delta n_{co} J^{(m)}}{\lambda} \tag{4}$$

And,  $J^{(m)}$  is the overlap integral between the fundamental guided mode and the cladding modes in the core area of the fiber [4].

In this research,  $CO_2$  laser-induced LPGs sensors of 500 µm pitch and 2.5 mm long imprinted on standard optical fibers were employed. The experimental setup is shown in figure 1. A super-luminescent LED is used as the light source. A photodetector coupled to an oscilloscope are used for taking measurements. The LPG is glued on a wooden bar, which has one fixed end and the other one is let free. The deflection of the bar (accounted as a vertical displacement  $\Delta y$  of the free bar end), causes the LPG to bend and, thus, its curvature response can be monitored.



Figure 1. (a) Scheme and (b) picture of the experimental setup.

When detecting light using the photodetector, as its voltage response is proportional to the overall optical power from the broadband light source, negative voltage variations are related to negative  $\Delta y$  values and positive voltage variations identify positive  $\Delta y$  values. The dependence of voltage on bar vertical displacement was seen to be linear [1].

# 3. 'Bar resonance' method results

Initially, the bar put in oscillation in a cantilever setup, is slightly deflected and then released. The signal measured in the oscilloscope takes a sinusoidal form whose amplitude decays as a function of time (figure 2a). By taking the Fourier transform of the measured signal, one can identify the frequency components of the oscillating bar movement. The figure 2b shows the amplitude of the fast Fourier transform (FFT), calculated from figure 2a data as a function of the frequency.



Figure 2. a) Photodetector time response of LPG sensor. b) The fast Fourier transform (FFT) amplitude as a function of frequency.

The value found for Kv (in all cases) shows that  $\left(\frac{k_v}{2\rho A}\right)^2 \ll \frac{YI}{\rho A} \left(\frac{\lambda_n}{l}\right)^4$ . The flexural elasticity modules value of the material of the cantilever (Y) is determined from the experimental data fifting frequency of the first normal mode (f<sub>1</sub>) versus the overhang length (L) of the cantilever.

Figure 3 shows the frequency setting of the first normal mode (f1) against the cantilevered length (L) for bars of five different woods (cachimbo, mahogany, cedar, pine and tornillo). The values determined for the flexural elasticity



moduli Y of the woods, based on the adjustment of equation (1) with their respective experimental data, are shown in table 2.

Figure 3. The first harmonic frequencies f1 (blue circles) versus the length of the cantilevered. The solid curve represents the fifting of equation (1) with the experimental data of the woods. The determined values of the flexural elasticity moduli of the woods are shown in each curve.

Table 1	shows the	values of	f the dime	ensions,	densities,	humidity	percentage	and moment	of inertia l	of the	e woods
studied.											

Table 1: Characteristics of the woods studied						
Material	h (m)	b (m)	ρ (kg/m <sup>3</sup> )	I (Kgm <sup>2</sup> ).10 <sup>-8</sup>	Moisture content (%)	
Cariniana domesticata Martius (Cachimbo) Swietenia mahagoni	0.0158	0.0513	590	1.6862	12,42	
(Caoba) Carapa guianensis Aubl.	0.0148	0.05155	469	1.3926	12,50	
(Cedro) Pinus patula	0.01615	0.05101	451	1.7906	11,59	
(Pino) Cedrelinga cateniformis	0.0149	0.04925	465	1.3576	10,91	
(Tornillo)	0.015/4	0.051/5	4/0	1.681/	11,39	

The flexural elasticity modules, determined by this technique, for the five woods studied, they were in good agreement with the order of magnitude reported for these modules, using different techniques [5-10].

Table 2: flexural elasticity moduli of the woods						
	Y	Y(reference)				
Material	(GPa)	(GPa)				
Cariniana domesticata Martius (Cachimbo)	11.53	9.29-12.84				
Swietenia mahagoni (Caoba)	7.05	7.2-10.3				
Carapa guianensis Aubl. (Cedro)	6.07	6.0-7.25				
Pinus patula		8.0-13.1				
(Pino)	10.1					
Cedrelinga cateniformis (Tornillo)	9.86	8.2-10.9				

#### Conclusions

This paper has reported the application of fiber long-period gratings to vibration monitoring and flexural elasticity modules determination. To the best of our knowledge, this is the first article to deal with the determination of wood material's flexural elasticity modules using long-period fiber gratings. Cachimbo, Mahagoni, Cedar, Pinus and Tornillo samples were tested. The values of the flexural elasticity moduli of the woods, determined by us, are in good agreement with the values reported in the literature using other techniques. This indicates that long-period fiber gratings can be straightforwardly employed in the dynamic characterization of a material's elastic properties.

# Acknowledgments

This research was supported by IGI-FIC and VRI of the National Engineering University. The authors thank for the support provided.

#### References

[1] L Mosquera, Jonas H Osório and Cristiano M B Cordeiro. "Determination of Young's modulus using optical fiber long-period gratings". Meas. Sci. Technol. 27 (2016).

[2] Rafael M. Digilov and Haim Abramovich. "Flexural Vibration Test of a Beam Elastically Restrained at One End: A New Approach for Young's Modulus Determination". Advances in Materials Science and Engineering (2013).

[3] S Tanaka, H Somatomo, A Wada, and N Takahashi. "Fiber-Optic Mechanical Vibration Sensor Using Long-Period Fiber Grating". Japanese Journal of Applied Physics 48 (2009).

[4] T. Erdogan and J. Lightwave, "Fiber Grating Spectra," Journal of Lightwave Technology, Vol. 15, No. 8, 1997, pp. 1277-1294.

[5] P Langbourd et all. "Comparison of wood properties of planted big-leaf Mahogany (Swietenia Macrophylla) in Martinique island with naturally grown Mahogany from Brasil, Mexico and Perú". Journal of tropical Forest Science 23(3):252-259 (2011).

[6] Yoza L., Baradit E. & Acevedo M., "Characterization of the physical mechanics properties of the peruvian species, pino (pinus patula) and tornillo (cedrelinga cateniformis) by using non destructives techniques". Anales Científicos, 76 (1): 12-16 (2015). DOI: http://dx.doi.org/10.21704/ac.v76i1.758

[7] JuliánAndrés Zárate Ramírez and Óscar Javier Gutiérrez Junco, "Wood species used in Boyacá characterization, according the Colombian Earthquake-Resistant Construction's Regulations". Engineering Faculty Journal, UPTC, January-June (2012), Vol. 21, No 32, pp.73-91.

[8] A. Aróstegui and A. Sato. "Study of the physical-mechanical properties of wood from 16 forest species in Peru". Revista Forestal del Perú V.4 (1-2):1-13.

[9] Eman Nabil et all. "Evaluation of Physical, Mechanical and Chemical Properties of Cedar and Sycamore Woods after Heat Treatment". Egypt. J. Chem. Vol. 61, No.6 pp. 1131 - 1149 (2018).

[10] Josefina S. Gonzalez. "Growth, properties and uses of western red cedar". Co-published by Western Red Cedar Lumber Association and Western Red Cedar Export Association. (Special publication, ISSN 0824-2119; no. SP-37R) (1997).